

# DETAILED STUDY MATERIALS: QUANTUM MECHANICS FUNDAMENTALS

## I. QUANTUM HARMONIC OSCILLATOR (QHO)

### A. INTRODUCTION AND SETUP

The Quantum Harmonic Oscillator (QHO) is a fundamental model in quantum mechanics, describing systems where a particle is subjected to a potential proportional to the square of its displacement from equilibrium. Examples include molecular vibrations and oscillations of atoms in a solid.

The potential energy is given by:

$$V(x) = 1/2 * m * \omega^2 * x^2$$

where:

- $m$  is the mass of the particle
- $\omega$  (omega) is the angular frequency of the oscillator
- $x$  is the displacement from equilibrium

The one-dimensional time-independent Schrödinger equation is:

$$- (\hbar^2 / 2m) * (d^2\psi(x) / dx^2) + V(x) * \psi(x) = E * \psi(x)$$

where:

- $\hbar$  (h-bar) is the reduced Planck's constant
- $\psi(x)$  (psi) is the wave function
- $E$  is the energy of the state

## B. SOLVING THE HERMITE DIFFERENTIAL EQUATION

Substituting  $V(x)$  into the Schrödinger equation and introducing dimensionless variables:

$$\xi = \sqrt{(m\omega/\hbar)} * x$$

The equation can be transformed into a form related to Hermite's differential equation:

$$(d^2u/d\xi^2) + ((2E / (\hbar\omega)) - \xi^2) * u = 0$$

where  $u(\xi)$  is proportional to  $\psi(x)$ . For bound states (wave functions that go to zero at infinity), solutions exist only for specific energy levels and involve Hermite polynomials,  $H_n(\xi)$ .

## C. ALGEBRAIC SOLUTION (LADDER OPERATORS)

An elegant method uses creation ( $a^\dagger$ ) and annihilation ( $a$ ) operators:

$$a = \sqrt{(m\omega/2\hbar)} * (x + i*p_x / (m\omega))$$

$$a^\dagger = \sqrt{(m\omega/2\hbar)} * (x - i*p_x / (m\omega))$$

with the commutation relation  $[a, a^\dagger] = 1$ .

The Hamiltonian operator can be expressed as:

$$x = \hbar\omega * (a^\dagger a + 1/2)$$

This leads to quantized energy levels:

$$E_n = (n + 1/2) * \hbar\omega$$

where  $n = 0, 1, 2, \dots$  is the principal quantum number.

- Ground state energy ( $n=0$ ):  $E_0 = 1/2 * \hbar\omega$
- Excited state energies:  $E_n = (n + 1/2) * \hbar\omega$

The corresponding wave functions are:

$$\psi_n(x) \propto H_n(\sqrt{(m\omega/\hbar)} * x) * \exp(-m\omega x^2/2\hbar)$$

## D. CLASSICAL TURNING POINTS

Classical turning points are the positions where the total energy  $E$  equals the potential energy  $V(x)$ . For the  $n$ -th energy level,  $E_n = V(x_{tp})$ :

$$(n + 1/2) * \hbar\omega = 1/2 * m * \omega^2 * x_{tp}^2$$

$$x_{tp} = \pm \sqrt{(2n + 1) * \hbar / (m\omega)}$$

These points define the classical limits of motion, within which the probability of finding the particle is non-zero. Quantum mechanically, there is a non-zero probability of finding the particle beyond these points (tunneling).

## E. EXPECTATION VALUES

For the  $n$ -th energy eigenstate  $\psi_n$ :

- Expectation value of position:  $\langle x \rangle_n = 0$
- Expectation value of position squared:  $\langle x^2 \rangle_n = ((n + 1/2) * \hbar) / (m\omega)$
- Expectation value of momentum:  $\langle p_x \rangle_n = 0$
- Expectation value of momentum squared:  $\langle p_x^2 \rangle_n = (m\omega/2) * (2n + 1) * \hbar$

## F. NUMERICAL PROBLEMS (EXAMPLES)

1. Calculate the energy levels for the first five states of a system with  $m = 1$  amu and  $\omega = 10^{13}$  rad/s. Determine the zero-point energy. [Hint: 1 amu  $\approx 1.66 \times 10^{-27}$  kg;  $\hbar \approx 1.054 \times 10^{-34}$  J\*s]

2. For the ground state ( $n=0$ ) of the QHO, find the classical turning points. If  $m=2$  kg and  $\omega=5$  rad/s, what is the energy? What are the turning points?

3. Using the ladder operators, show that  $[x, a] = -\hbar\omega a$ . (This demonstrates that 'a' lowers the energy state).

## II. RIGID ROTATOR

### A. ANGULAR MOMENTUM COMMUTATION RULES

Angular momentum is quantized. The components of the angular momentum operator ( $L_x, L_y, L_z$ ) do not commute:

$$[L_x, L_y] = i\hbar L_z$$

$$[L_y, L_z] = i\hbar L_x$$

$$[L_z, L_x] = i\hbar L_y$$

However, the square of the total angular momentum,  $L^2 = L_x^2 + L_y^2 + L_z^2$ , commutes with each component:

$$[L^2, L_i] = 0 \text{ for } i = x, y, z.$$

### B. ANGULAR MOMENTUM OPERATORS IN SPHERICAL POLAR COORDINATES

In spherical coordinates  $(r, \theta, \phi)$ , the angular momentum operators are typically expressed using derivatives with respect to angles:

$$L_z = -i\hbar (\partial/\partial\phi)$$

$$L^2 = -\hbar^2 [ 1/\sin(\theta) * \partial/\partial\theta * (\sin(\theta) * \partial/\partial\theta) + 1/\sin^2(\theta) * (\partial^2/\partial\phi^2) ]$$

### C. QUANTIZATION OF ANGULAR MOMENTUM

The simultaneous eigenstates of  $L^2$  and  $L_z$  are characterized by quantum numbers  $l$  and  $m$ :

- Eigenvalues of  $L^2$ :  $L^2 |l, m\rangle = l(l+1)\hbar^2 |l, m\rangle$ , where  $l = 0, 1, 2, \dots$
- Eigenvalues of  $L_z$ :  $L_z |l, m\rangle = m\hbar |l, m\rangle$ , where  $m = -l, -l+1, \dots, l-1, l$

Thus, the magnitude of the angular momentum vector is  $\sqrt{l(l+1)}\hbar$ , and its projection onto the z-axis is  $m\hbar$ .

## D. RIGID ROTATOR MODEL

A rigid rotator consists of two masses ( $m_1, m_2$ ) connected by a rigid rod of length  $r$ . For a diatomic molecule, this is approximated by reducing the system to a single mass  $\mu$  (reduced mass) at a fixed distance  $R$  from the origin.

$$\mu = (m_1 * m_2) / (m_1 + m_2)$$

The potential energy is zero ( $V=0$ ) as the bond length is fixed.

## E. SCHRÖDINGER EQUATION FOR RIGID ROTATOR

The Hamiltonian is simply the kinetic energy operator in terms of angular momentum:

$$x = L^2 / (2I)$$

where  $I = \mu R^2$  is the moment of inertia.

The Schrödinger equation is:

$$(L^2 / (2I)) * \psi(\theta, \phi) = E * \psi(\theta, \phi)$$

In spherical coordinates (where  $r$  is fixed):

$$- (\hbar^2 / 2I) * [ 1/\sin(\theta) * \partial/\partial\theta * (\sin(\theta) * \partial/\partial\theta) + 1/\sin^2(\theta) * \partial^2/\partial\phi^2 ] * \psi(\theta, \phi) = E * \psi(\theta, \phi)$$

## F. SEPARATION OF VARIABLES AND SPHERICAL HARMONICS

Separating the wave function into angular parts:  $\psi(\theta, \phi) = Y_{lm}(\theta, \phi)$ . The solutions are the Spherical Harmonics,  $Y_{lm}$ , which are simultaneous eigenfunctions of  $L^2$  and  $L_z$ .

The energy levels are:

$$E_l = l(l+1) * \hbar^2 / (2I)$$

The energy depends only on the quantum number  $l$  (the azimuthal quantum number), not on  $m$ . This leads to degeneracy.

## G. DISCUSSION OF SOLUTION

The Spherical Harmonics,  $Y_{lm}(\theta, \phi)$ , are complex functions that describe the angular distribution of the wave function. The solution implies that the rotational energy levels of a diatomic molecule are quantized, with the spacing between levels increasing with  $l$ .

## H. NUMERICAL PROBLEMS (EXAMPLES)

1. For a CO molecule (reduced mass  $\mu = 1.14 \times 10^{-26}$  kg, bond length  $R = 1.13 \times 10^{-10}$  m), calculate the moment of inertia. Then, find the energies for the first three rotational levels ( $l=0, 1, 2$ ).
2. Calculate the frequency of photons that would be emitted or absorbed during transitions between the  $l=1$  and  $l=2$  rotational states of the CO molecule. [Hint:  $E = \hbar\nu$ ]

## III. HYDROGEN ATOM AND HYDROGEN-LIKE IONS

## A. SETUP OF SCHRÖDINGER EQUATION

For a hydrogen atom (one electron of mass  $m$  and charge  $-e$  orbiting a nucleus of charge  $+Ze$ ), the potential energy is the Coulomb potential:

$$V(r) = - (Ze^2) / (4\pi\epsilon_0 r)$$

where  $\epsilon_0$  is the permittivity of free space. The Schrödinger equation is solved in spherical polar coordinates  $(r, \theta, \phi)$ , with the Hamiltonian being:

$$H = - (\hbar^2 / 2\mu) \nabla^2 + V(r)$$

where  $\mu$  is the reduced mass of the electron-nucleus system (for hydrogen, approximately the electron mass).

## B. SEPARATION OF VARIABLES

The wave function  $\psi(r, \theta, \phi)$  is separated into radial and angular parts:

$$\psi(r, \theta, \phi) = R(r) * Y_{lm}(\theta, \phi)$$

where  $R(r)$  is the radial wave function and  $Y_{lm}(\theta, \phi)$  are the Spherical Harmonics solved for the rigid rotator.

This separation leads to two equations: one for the radial part and one for the angular part (which yields the spherical harmonics and their associated quantum numbers  $l$  and  $m$ ).

## C. SOLUTION OF ANGULAR PART ( $\phi$ PART ONLY)

The part of the angular solution dependent on  $\phi$  arises from the  $L_z$  operator. The quantization condition for  $L_z$  gives:

$$L_z Y_{lm} = m\hbar Y_{lm}$$

where  $m$  is the magnetic quantum number ( $m = -l, \dots, l$ ). The solution for the  $\phi$  dependence is typically of the form  $e^{im\phi}$ .

## D. QUANTIZATION OF ENERGY

Solving the radial equation yields the quantized energy levels for the hydrogen atom ( $Z=1$ ):

$$E_n = -(\mu * (Ze^2)^2) / (2 * (4\pi\epsilon_0)^2 * n^2) = - (13.6 \text{ eV} * Z^2) / n^2$$

where  $n$  is the principal quantum number ( $n = 1, 2, 3, \dots$ ). The energy depends only on  $n$ , leading to degeneracy.

## E. REAL WAVE FUNCTIONS

The spherical harmonics  $Y_{lm}$  are complex. For some applications, real wave functions are preferred. These can be constructed as linear combinations of complex  $Y_{lm}$  functions with the same / but different  $m$  values (e.g.,  $Y_{l,m} + Y_{l,-m}$  and  $i(Y_{l,m} - Y_{l,-m})$ ).

## F. AVERAGE AND MOST PROBABLE DISTANCES

The radial wave function  $R_{nl}(r)$  determines the probability distribution of the electron's distance from the nucleus.

- Most Probable Distance ( $r_{mp}$ ): Found by maximizing the radial probability density  $P(r) = r^2 |R_{nl}(r)|^2$ .
- Average Distance ( $\langle r \rangle_{nl}$ ): Calculated by integrating  $r * P(r)$  over all  $r$ .

For Hydrogen ( $Z=1$ ):

- $\langle r \rangle_{nl} = (a_0/2) * [3n^2 - l(l+1)]$ , where  $a_0 = 4\pi\epsilon_0 \hbar^2 / (m_e e^2)$  is the Bohr radius.
- Most probable distance is also related to  $n$  and  $l$ , specific formulas vary. For the ground state ( $n=1, l=0$ ),  $\langle r \rangle_{10} = 1.5 a_0$ , and  $r_{mp} = a_0$ .

## G. SCHRÖDINGER EQUATION FOR MANY-ELECTRON ATOMS

Setting up the Schrödinger equation for atoms with more than one electron (like Helium,  $Z=2$ , with 2 electrons) becomes significantly more complex due to the electron-electron repulsion term in the potential:

$$V_{\text{total}} = V_{\text{nucleus}} + \text{Sum}(V_{\text{electron\_nucleus}}) + \text{Sum}(V_{\text{electron\_electron}})$$

For Helium, the interaction term is  $e^2 / (4\pi\epsilon_0 |r_1 - r_2|)$ , which depends on the distance between the two electrons ( $r_1$  and  $r_2$ ).

This term makes the equation non-separable in spherical coordinates, requiring approximations (like the Hartree-Fock method or variational methods) to find solutions.

## H. NUMERICAL PROBLEMS (EXAMPLES)

1. Calculate the ionization energy of the ground state of  $\text{He}^+$  ion ( $Z=2$ ). [Hint: Use the energy formula for hydrogen-like ions and note ionization means removing the electron to  $n=\infty$ ]
2. For the Hydrogen atom's ground state ( $n=1, l=0$ ), calculate the most probable distance and the average distance of the electron from the nucleus in terms of the Bohr radius ( $a_0$ ).
3. Explain why the Schrödinger equation for Helium cannot be solved by simple separation of variables.